

High Order Well-Balanced Discrete Kinetic Model for Shallow Water Equations

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joint work with Mario Ricchiuto and R emi Abgrall

- 1 Models
- 2 Residual Distribution
- 3 Time Discretization
 - IMEX
 - Deferred Correction
- 4 Structure preserving
- 5 Numerical tests
- 6 Conclusion and perspective

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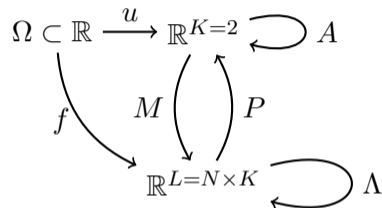
Kinetic and Shallow water equations

Modify the kinetic relaxation models by D. Aregba-Driollet and R. Natalini¹

Hyperbolic limit equation is

$$(1) \quad u_t + \partial_x A(u) = 0, \quad u : \Omega \rightarrow \mathbb{R}^K$$

(2)



Relaxation system

$$(3) \quad f^\varepsilon = (f_1, f_2, \dots, f_N) = (h_1, q_1, h_2, q_2, \dots, h_N, q_N)$$

$$(4) \quad f_t^\varepsilon + \Lambda \partial_x f^\varepsilon = \frac{1}{\varepsilon} (M(P f^\varepsilon) - f^\varepsilon), \quad f^\varepsilon : \Omega \rightarrow \mathbb{R}^L$$

$$P f^\varepsilon \rightarrow u, \quad P(M(u)) = u, \quad P \Lambda M(u) = A(u)$$

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Kinetic and Shallow water equations

Modify the kinetic relaxation models by D. Aregba-Driollet and R. Natalini¹

Hyperbolic limit equation is

$$(1) \quad u_t + \partial_x A(u) + S_b(u) = 0, \quad u : \Omega \rightarrow \mathbb{R}^K$$

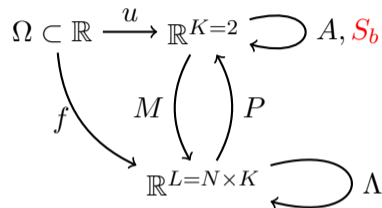
$$(2) \quad \begin{cases} h_t + q_x = 0 \\ q_t + (q^2/h + \frac{g}{2}h^2)_x + ghb_x = 0 \end{cases}$$

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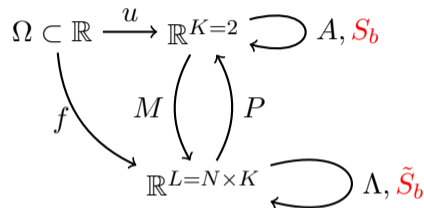
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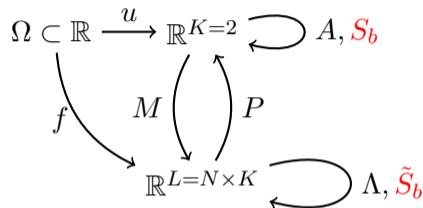
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$$P(M(u)) = u, \quad P\Lambda M(u) = A(u), \quad P\tilde{S}_b(f) = S_b(Pf), \quad P\Lambda\tilde{S}_b(f) = S_b(P\Lambda f).$$



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$$f^\varepsilon = (f_1, f_2)^T = (h_1, q_1, h_2, q_2)^T$$

Diagonal relaxation method (DRM)

- $K = 2$
- $N = D + 1 = 2$
- $L = N \times K = 2 \times 2$
- $P = (I_K, \dots, I_K) = (I_2, I_2)$
- $\Lambda = \begin{pmatrix} -\lambda I_2 & \\ & \lambda I_2 \end{pmatrix}$
- $M_1(u) = \frac{u\lambda - A(u)}{2\lambda}$
- $M_2(u) = \frac{u\lambda + A(u)}{2\lambda}$

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$$u^\varepsilon := P f^\varepsilon = f_1 + f_2$$

$$h^\varepsilon = h_1 + h_2, \quad q^\varepsilon = q_1 + q_2$$

$$f_t^\varepsilon + \Lambda \partial_x f^\varepsilon + \tilde{S}_b(f^\varepsilon) = \frac{1}{\varepsilon} (M(P f^\varepsilon) - f^\varepsilon)$$

$$\partial_t \begin{pmatrix} h_1 \\ q_1 \\ h_2 \\ q_2 \end{pmatrix} + \partial_x \begin{pmatrix} -\lambda h_1 \\ -\lambda q_1 \\ \lambda h_2 \\ \lambda q_2 \end{pmatrix} - \begin{pmatrix} 0 \\ g(h_1 + \frac{b}{2}) \partial_x b \\ 0 \\ g(h_2 + \frac{b}{2}) \partial_x b \end{pmatrix} =$$

$$\frac{1}{2\varepsilon} \begin{pmatrix} -h_1 + h_2 - \frac{q^\varepsilon}{\lambda} \\ -q_1 + q_2 - \frac{(q^\varepsilon)^2 / (h^\varepsilon) + g((h^\varepsilon)^2 - b^2) / 2}{\lambda} \\ h_1 - h_2 + \frac{q^\varepsilon}{\lambda} \\ q_1 - q_2 + \frac{(q^\varepsilon)^2 / (h^\varepsilon) + g((h^\varepsilon)^2 - b^2) / 2}{\lambda} \end{pmatrix}.$$

Relaxation system

$$f_t^\varepsilon + \Lambda \partial_x f^\varepsilon + \tilde{S}_b(f^\varepsilon) = \frac{M(Pf^\varepsilon) - f^\varepsilon}{\varepsilon},$$
$$+$$
$$P(M(u)) = u, \quad P\Lambda M(u) = A(u),$$
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$$\partial_t u^\varepsilon + \partial_x A(u^\varepsilon) + S_b(u^\varepsilon) = \varepsilon \Xi + \mathcal{O}(\varepsilon^2), \quad u^\varepsilon = Pf^\varepsilon,$$

where $\Xi := \partial_x(B(u^\varepsilon)\partial_x u^\varepsilon) + \partial_x(-A'(u^\varepsilon)S_b(u^\varepsilon) + S_b(A(u^\varepsilon)))$,

with $B(u) := P\Lambda^2 M'(u) - A'(u)^2 \in \mathbb{R}^{S \times S}$.

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Whitham's subcharacteristic condition $B \geq 0 \implies$ Diffusive

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- High order
- FE based
- Compact stencil
- No need of conservative variables
- Can recast some other FV, FE schemes²

Finite Element Setting

$$\partial_t f + \nabla_x \cdot A(f) = S(f)$$

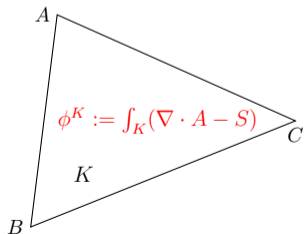
$$V_h = \{f \in L^2(\Omega_h, \mathbb{R}^L) \cap \mathcal{C}^0(\Omega_h), \\ f|_K \in \mathbb{P}^p, \forall K \in \Omega_h\}$$

$$f(x) = \sum_{\sigma \in D_h} f_\sigma \varphi_\sigma(x) \\ = \sum_{K \in \Omega_h} \sum_{\sigma \in K} f_\sigma \varphi_\sigma(x)|_K$$

²R. Abgrall. Some remarks about conservation for residual distribution schemes. Computational Methods in Applied Mathematics, 2018. DOI: <https://doi.org/10.1515/cmam-2017-0056>.

Residual Distribution - Spatial Discretization

- 1 Define $\forall K \in \Omega_h$ a fluctuation term (total residual) $\phi^K = \int_K \nabla \cdot A(f) - S(f) dx$

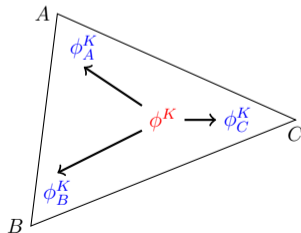
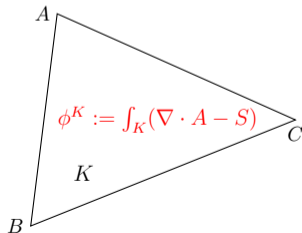


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- 2 Define nodal residuals $\phi_\sigma^K \forall \sigma \in K : \phi^K = \sum_{\sigma \in K} \phi_\sigma^K, \quad \forall K \in \Omega_h.$

Choice of Residuals

Basic algorithm (Galerkin), numerical fluxes (Rusanov), linear stabilization terms (SUPG, jump derivative penalty), non linear stabilization (PSI).



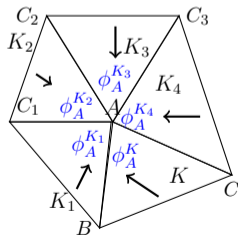
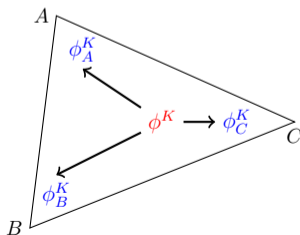
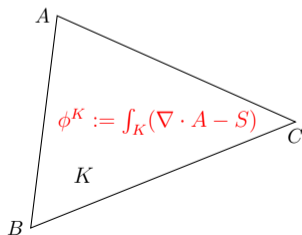
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Choice of Residuals

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- 3 The resulting scheme is $\partial_t f_\sigma + \sum_{K|\sigma \in K} \phi_\sigma^K = 0, \quad \forall \sigma \in D_h.$



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Stiff source term \Rightarrow unstable for $\varepsilon \ll \Delta t \Rightarrow$ IMEX approach:

IMplicit for stiff source term, EXplicit for advection term and bathymetry source

$$(5) \quad \frac{f^{n+1,\varepsilon} - f^{n,\varepsilon}}{\Delta t} + \Lambda \partial_x f^{n,\varepsilon} + \tilde{S}_b(f^{n,\varepsilon}) = \frac{1}{\varepsilon} (M(P f^{n+1,\varepsilon}) - f^{n+1,\varepsilon}).$$

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How to treat non-linear implicit functions?

Recall: $PM(u) = u$ and $Pf^\varepsilon = u^\varepsilon$, so

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Find $u^{n+1,\varepsilon} = Pf^{n+1,\varepsilon}$ and substitute it in the Maxwellian in (5).

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- IMEX formulation is first order accurate $=: \mathcal{L}^1$
- IMEX formulation is asymptotic preserving (AP) (as $\varepsilon \rightarrow 0$ we recast SW)

Deferred Correction³

How to combine two methods keeping the accuracy of the second and the stability and simplicity of the first one?

$$f^{0,(k)} := f(t^n), \quad k = 0, \dots, K,$$

$$f^{m,(0)} := f(t^n), \quad m = 1, \dots, M$$

$$\mathcal{L}^1(f^{(k)}) = \mathcal{L}^1(f^{(k-1)}) - \mathcal{L}^2(f^{(k-1)}) \text{ with } k = 1, \dots, K.$$

DeC Theorem

- \mathcal{L}^1 coercive
- $\mathcal{L}^1 - \mathcal{L}^2$ Lipschitz

DeC converges and $\min(K, M + 1)$ is the order of accuracy.

$$\mathcal{L}^1(f) = 0$$

- IMEX
- First order accurate
- Mass lumping
- Computationally explicit

$$\mathcal{L}^2(f) = 0$$

- Order $M + 1$
- Quadrature in timestep
- Nonlinearly implicit
- Implicit Runge–Kutta

³A. Dutt, L. Greengard, and V. Rokhlin. BIT Numerical Mathematics, 40(2):241–266, 2000.

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Other properties

- Well balancedness: lake at rest steady state preservation

- Match of the discretizations of the source term and the flux when $v = 0$ and

$$\eta(x) = h^\varepsilon(x) + b(x) \equiv \eta_0$$

- $\phi_\sigma^K = \int_K g\varphi^\sigma \partial_x \frac{(h^\varepsilon)^2 - b^2}{2} dx + \int_E g\varphi^\sigma (h^\varepsilon + b) \partial_x b dx = 0$

$$\int_K g\varphi^\sigma \partial_x \varphi^i(x) \underbrace{\frac{h^\varepsilon(x_i) - b(x_i)}{2}}_{=\frac{\eta_0}{2} - b(x_i)} \underbrace{(h^\varepsilon(x_i) + b(x_i))}_{=\eta_0} dx = \int_K -g\varphi^\sigma \eta_0 \partial_x \varphi^i(x) b(x_i) dx =$$
$$- \int_E g\varphi^\sigma (h^\varepsilon + b) \partial_x b dx.$$

- Recipe for all sources \tilde{S}_b
- Stabilization techniques depends on η instead of h

- Depth non-negativity
 - Wet and dry elements: dry elements $h_1 + h_2 \leq \tau_{010}$, $q_1 = q_2 = 0$
 - Hybrid elements \implies Lower the bathymetry to have positive DoFs where $h_1 + h_2 \leq \tau_{010}$ (IC or along the scheme) to preserve the well-balancedness $\eta_K \equiv C$
 - Use of explicit schemes that, subjected to CFL conditions, can preserve the non-negativity, e.g. Rusanov, modified Rusanov, blended PSI+Rusanov. (Not jump stabilization)

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Simulations: Convergence for Subcritical Flow

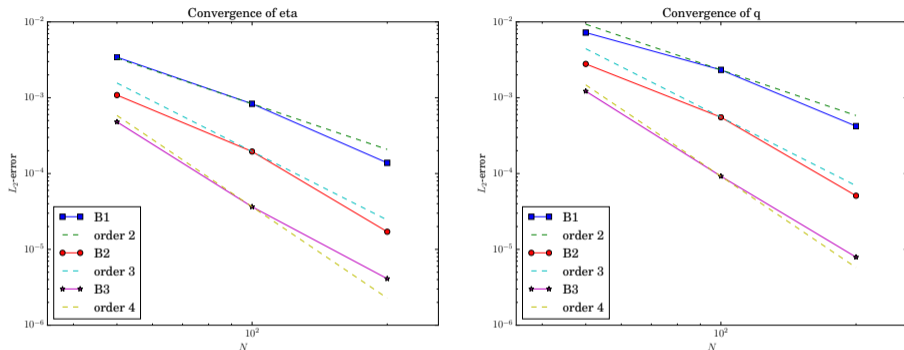


Figure: Subcritical flow: convergence for $\eta^\varepsilon = h^\varepsilon + b$ and $h^\varepsilon v^\varepsilon$

$$b(x) = \begin{cases} 0.2 \exp\left(\frac{((x-10)/5)^2}{1-((x-10)/5)^2}\right), & \text{if } x \in B_5(10), \\ 0, & \text{else.} \end{cases} \quad \begin{aligned} h^\varepsilon(0, x) &= 2 - b(x) & q^\varepsilon(0, t) &= 4.42 \\ q^\varepsilon(0, x) &= 4.42 & h^\varepsilon(25, t) &= 2 \end{aligned}$$

$$\lambda = 6.5, \quad \varepsilon = 10^{-14}, \quad \begin{aligned} f^\varepsilon(0, x) &= M(u^\varepsilon(0, x)) & T &= 100 \end{aligned}$$

Simulation: transcritical with shock

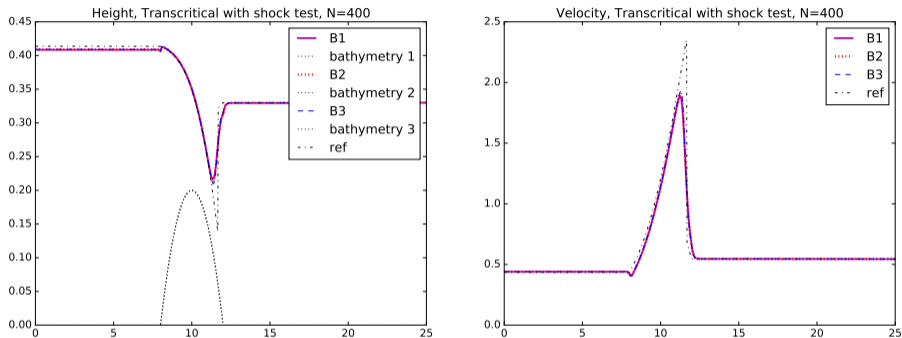


Figure: Transcritical flow with shock test: η^ϵ and v^ϵ with $N = 400$

$$b(x) = (0.2 - 0.05(x - 10)^2) \mathbb{1}_{\{8 < x < 12\}}$$

$$\eta^\epsilon(0, x) = 0.4 - 0.07 \mathbb{1}_{\{x > 8\}}$$

$$q^\epsilon(0, x) = 0.14$$

$$q^\epsilon(0, t) = 0.18$$

$$h^\epsilon(25, t) = 0.33$$

$$\lambda = 4, \quad \epsilon = 10^{-14}.$$

Simulations: lake at rest

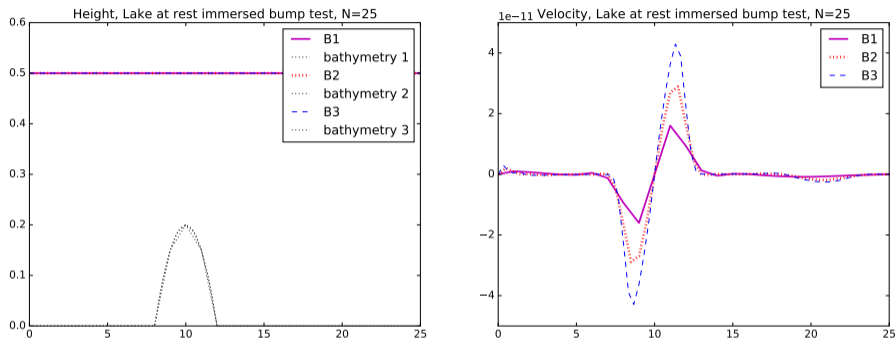


Figure: Lake at rest with immersed bump test: η^ε and v^ε with $N = 25$

$$b(x) = (0.2 - 0.05(x - 10)^2)\mathbb{1}_{\{8 < x < 12\}}$$

$$\eta^\varepsilon(0, x) = 0.5$$

$$q^\varepsilon(0, x) = 0$$

$$q^\varepsilon(0, t) = 0$$

$$q^\varepsilon(25, t) = 0$$

$$\lambda = 2$$

$$q - q^{ex} = \mathcal{O}(N_t \varepsilon)$$

$$T = 3$$

$$\varepsilon = 10^{-14}$$

Simulations: wet and dry lake at rest

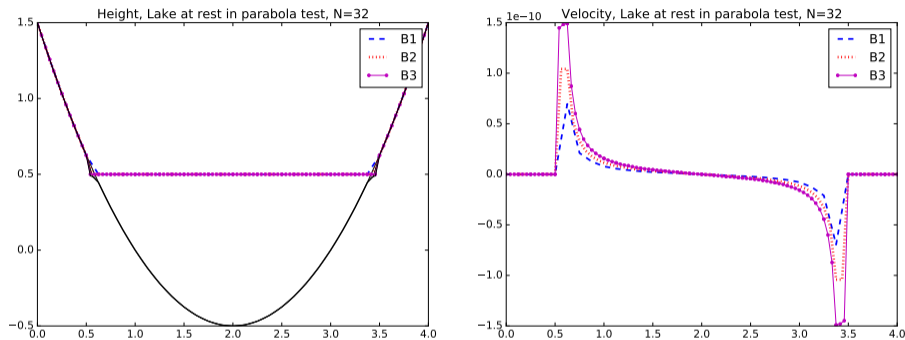
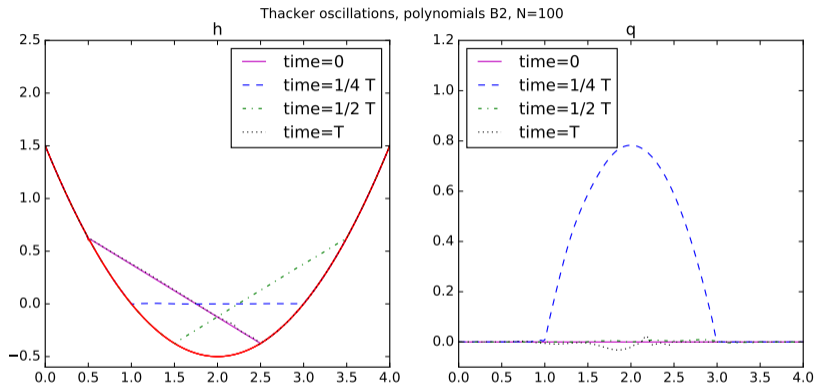


Figure: Lake at rest in parabola test: η^ε and v^ε with $N = 32$

$$b(x) = (x - 2)^2 - 0.5$$
$$\eta^\varepsilon(0, x) = \max(0.5, b(x))$$
$$\lambda = 4$$

$$q - q^{ex} = \mathcal{O}(N_t \varepsilon)$$
$$T = 3$$
$$\varepsilon = 10^{-14}$$

Simulations: Thicker Oscillations



$$b(x) = (x - 2)^2 - 0.5$$

$$\eta^\varepsilon(0, x) = \max(-0.5x + 0.875, b(x))$$

$$\lambda = 6.5$$

$$\text{period} = 2.0606$$

$$T = 5 \cdot 2.0606$$

$$\varepsilon = 10^{-14}$$

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Conclusions

- IMEX
- High order
- Residual Distribution
- Deferred Correction
- Well-balanced
- Wet/dry
- Nonnegative water height

Perspective

- MOOD
- Entropy stability
- Multi dimension

- 1 R. Abgrall, and D.T.. High Order Asymptotic Preserving Deferred Correction Implicit-Explicit Schemes for Kinetic Models. *SIAM Journal on Scientific Computing*, 42(3):B816–B845, 2020.
- 2 D. Aregba-Driollet, and R. Natalini. Discrete kinetic schemes for multidimensional systems of conservation laws. *SIAM J. Numer. Anal.*, 37(6):1973–2004, 2000.
- 3 A. Dutt, L. Greengard, and V. Rokhlin. Spectral Deferred Correction Methods for Ordinary Differential Equations. *BIT Numerical Mathematics*, 40(2):241–266, 2000.
- 4 R. Abgrall. High Order Schemes for Hyperbolic Problems Using Globally Continuous Approximation and Avoiding Mass Matrices. *Journal of Scientific Computing*, 73(2):461–494, 2017.
- 5 M. Ricchiuto, and A. Bollermann. Stabilized residual distribution for shallow water simulations. *Journal of Computational Physics*, 228(4):1071–1115, 2009.

Thank you for the attention!